

Dynamical large deviations of Markov processes

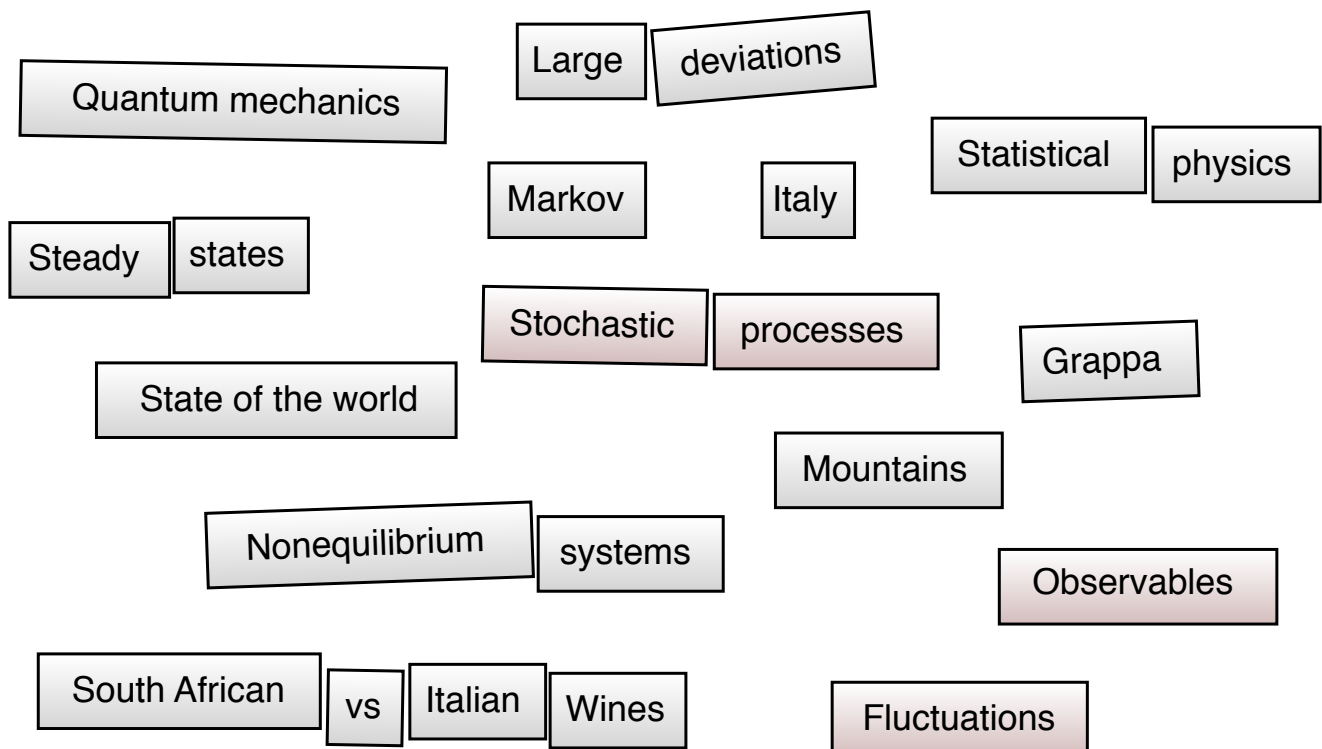
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Lecture notes: [arxiv:1705.06492](https://arxiv.org/abs/1705.06492)

Detailed plan



Bar questions

Question 1

Is there,
can there be,
will there be
a complete theory of nonequilibrium systems?

Question 2

What is nonequilibrium?

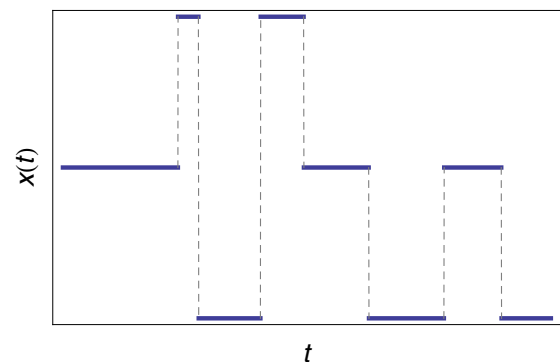
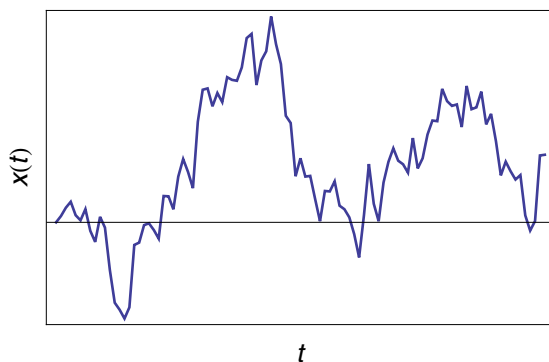
Question 3

Why is nonequilibrium more difficult than equilibrium?

Question 4

Where's the bar?

Langevin equations



- SDE:

$$dX_t = F(X_t)dt + \sigma dW_t, \quad X_t \in \mathbb{R}^n, W_t \in \mathbb{R}^m$$

- Drift: $F(x) \in \mathbb{R}^d$
- Noise matrix: σ ($n \times m$)

- Noisy ODE:

$$\dot{X}_t = F(X_t) + \sigma \xi_t$$

- White noise:

$$\langle \xi_t \rangle = 0, \quad \langle \xi_t^i \xi_{t'}^j \rangle = \delta_{ij} \delta(t - t'),$$

Fokker-Planck equation

- Fokker-Planck equation:

$$\partial_t p(x, t) = L^\dagger p(x, t), \quad p(x, t) = P(X_t = x)$$

- Fokker-Planck generator:

$$L^\dagger = -\nabla \cdot F + \frac{1}{2} \nabla \cdot D \nabla$$

- Conservation equation:

$$\partial_t p(x, t) + \nabla \cdot J_t(x) = 0$$

- Fokker-Planck current:

$$J_t(x) = F(x)p(x, t) - \frac{D}{2} \nabla p(x, t)$$

- Stationary distribution:

$$L^\dagger p_s = 0$$

Examples

Kramers or underdamped Langevin equation

$$\begin{aligned} dq_t &= \frac{p_t}{m} dt \\ dp_t &= \left(-\nabla V(q_t) + \phi_t - \Gamma \frac{p_t}{m} \right) dt + \sqrt{2\Gamma/\beta} dW_t \end{aligned}$$

Overdamped Langevin equation

$$dq_t = \Gamma^{-1} (-\nabla V + \phi_t) dt + \sqrt{2\Gamma^{-1}/\beta} dW_t.$$

Gradient SDEs

$$dX_t = -\nabla U(X_t) dt + \sigma dW_t$$

Stationary distribution:

$$p_s(x) = c e^{-2U(x)/\varepsilon^2}$$

Linear diffusions

$$dX_t = -MX_t dt + \sigma dW_t$$

- Stationary distribution:

$$p_s(x) = \sqrt{\frac{\det C}{2\pi}} \exp\left(-\frac{1}{2}x \cdot C^{-1}x\right)$$

- Lyapunov equation:

$$D = MC + CM^T, \quad D = \sigma\sigma^T$$

- Ornstein-Uhlenbeck process:

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

Equilibrium vs nonequilibrium

Equilibrium

- Example: Gradient SDE
- Detailed balance
- Reversible process
- $J_s(x) = 0$
- Spectrum of L real
- Has no “rotation”

Nonequilibrium

- Example: Non-gradient SDE
- No detailed balance
- Nonreversible process
- $J_s(x) \neq 0$ but $\nabla \cdot J_s = 0$
- Spectrum of L complex
- Has “rotation”

Question 5

Why study Markov processes?

Examples

- Linear process:

$$dX_t = -MX_t dt + \sigma dW_t$$

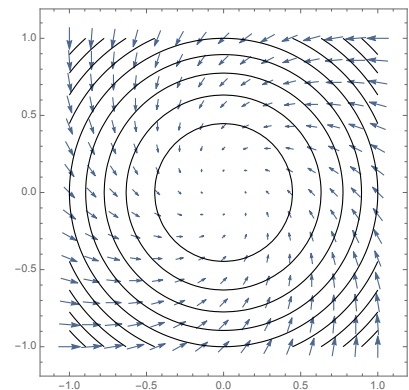
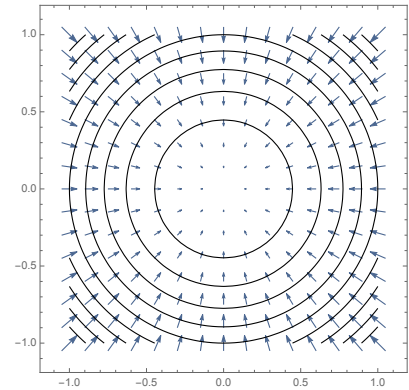
- Noise matrix: $\sigma = \epsilon \mathbb{1}$
- Gradient system: $M = \mathbb{1}$

$$F = -\nabla U, \quad U(x) = \frac{x^2}{2}$$

- Non-gradient system:

$$M = \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$$

- Depends in general on M and $D = \sigma\sigma^T$



Evolution of expectations

- Average or expectation:

$$\langle f(X_t) \rangle = E[f(X_t)] = \int f(x) p(x, t) dx$$

- Evolution:

$$\partial_t \langle f(X_t) \rangle = \langle (Lf)(X_t) \rangle$$

- Generator:

$$L = (L^\dagger)^\dagger = F \cdot \nabla + \frac{1}{2} \nabla \cdot D \nabla$$

- $L \neq L^\dagger$ if $F \neq 0$

Scrödinger picture

$$\partial_t p(x, t) = L^\dagger p(x, t)$$

Heisenberg picture

$$\partial_t \langle f(X_t) \rangle = \langle (Lf)(X_t) \rangle$$

Non-Hermitian operators

- Inner or scalar product:

$$\langle f(X) \rangle = \int p(x) f(x) dx = \langle p, f \rangle$$

- Duality or adjoint:

$$\langle p, Lf \rangle = \langle L^\dagger p, f \rangle$$

- Fokker-Planck equation:

$$\partial_t \langle f(X_t) \rangle = \langle p_t, Lf \rangle = \langle L^\dagger p_t, f \rangle$$

- Differential operator:

$$\langle p, \nabla f \rangle = \int p \nabla f dx = \int p df = \underbrace{pf|_{\text{boundary}}}_{=0} - \int f dp$$

- $\nabla^\dagger = -\nabla$ (skew symmetric)
- $\Delta^\dagger = \Delta$

Non-Hermitian operators (cont'd)

- Direct problem:

$$Lv(x) = \lambda v(x)$$

- Dual problem:

$$L^\dagger u(x) = \beta u(x), \quad \beta = \lambda^*$$

- Orthonormal basis:

$$\langle u_i, v_j \rangle = \int u_i^*(x) v_j(x) dx = \delta_{ij}$$

- Completeness:

$$\delta(x - x') = \sum_i u_i^*(x) v_i(x')$$

Comparison with quantum mechanics

	Markov	Quantum
State	X_t	$ \psi(t)\rangle$ or $\psi(x, t)$
Distribution	$p(x, t)$	$ \psi(x, t) ^2$
Evolution	Fokker-Planck	Schrödinger
Generator	L	H
Propagator	$U(t) = e^{L^\dagger t}$	$U(t) = e^{-iHt/\hbar}$
Inner product	$\langle p, f \rangle$	$\langle \psi, \psi \rangle = \langle \psi \psi \rangle$
Duality	$\langle p, Af \rangle = \langle A^\dagger p, f \rangle$	$\langle \psi, A\psi \rangle = \langle A\psi, \psi \rangle = \langle \psi A \psi \rangle$
Self-adjoint?	Not necessarily	Always (closed systems)

Dynamical observables

- General observable

$$A_T = \frac{1}{T} \int_0^T f(X_t) dt + \frac{1}{T} \int_0^T g(X_t) \circ dX_t$$

- Increment: $dX_t = X_{t+dt} - X_t$
- Stratonovich product:

$$g(X_t) \circ dX_t = g(X_t + dX_t/2) dX_t = \underbrace{g(X_t) dX_t}_{\text{Ito}} + \frac{g'(X_t)}{2} dt$$

- Interpretation:

$$A_T = \underbrace{\frac{1}{T} \int_0^T f(X_t) dt}_{\text{state}} + \underbrace{\frac{1}{T} \int_0^T g(X_t) \dot{X}_t dt}_{\text{velocity}}$$

Examples

- Potential energy

$$\Delta U_T = U(X_T) - U(X_0) = \int_0^T \nabla U(X_t) \circ dX_t$$

- Work:

$$W_T = \int_0^T F(X_t) \circ dX_t$$

- Entropy production:

$$\Sigma_T = 2 \int_0^T (D^{-1}F(X_t)) \circ dX_t$$

- Empirical density:

$$\rho_T(x) = \frac{1}{T} \int_0^T \delta(X_t - x) dt \rightarrow p_s(x)$$

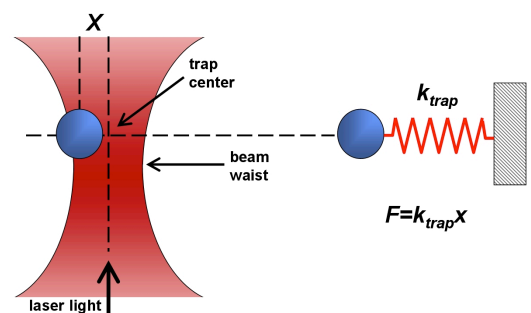
- Empirical current:

$$J_T(x) = \frac{1}{T} \int_0^T \delta(X_t - x) \circ dX_t \rightarrow J_s(x)$$

Example: Pulled Brownian particle

- Glass bead in water
- Laser tweezers
- Langevin dynamics:

$$m\ddot{x}(t) = \underbrace{-\alpha\dot{x}}_{\text{drag}} - \underbrace{k[x(t) - vt]}_{\text{spring force}} + \underbrace{\xi(t)}_{\text{noise}}$$

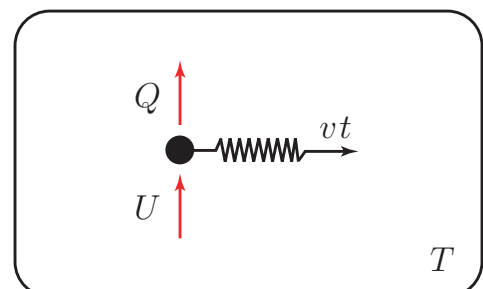


- Fluctuating work:

$$\underbrace{W_T}_{\text{work}} = \underbrace{\Delta U}_{\text{potential}} + \underbrace{Q_T}_{\text{heat}}$$

- Work observable:

$$W_T = -kv \int_0^T [x(t) - vt] dt$$



- Work distribution: $P(W_T = w)$

Large deviation theory

- Random variable: A_T
- Probability density: $P(A_T = a)$

Large deviation principle (LDP)

$$P(A_T = a) \approx e^{-TI(a)}$$

- Meaning of \approx :

$$\ln P(a) = -TI(a) + o(T)$$
$$\lim_{T \rightarrow \infty} -\frac{1}{T} \ln P(a) = I(a)$$

- Rate function: $I(a) \geq 0$

Goals of large deviation theory

- 1 Prove large deviation principle exists
- 2 Calculate rate function

Varadhan's Theorem

- LDP:

$$P(A_T = a) \approx e^{-TI(a)}$$

- Exponential expectation:

$$\langle e^{Tf(A_T)} \rangle = \int e^{Tf(a)} P(A_T = a) da$$

- Limit functional:

$$\lambda(f) = \lim_{T \rightarrow \infty} \frac{1}{T} \ln \langle e^{Tf(A_T)} \rangle$$



S. R. Srinivasa Varadhan
Abel Prize 2007

Theorem: Varadhan (1966)

$$\lambda(f) = \max_a \{f(a) - I(a)\}$$

Special case: $f(a) = ka$

$$\lambda(k) = \max_a \{ka - I(a)\}$$

Gärtner-Ellis Theorem

Scaled cumulant generating function (SCGF)

$$\lambda(k) = \lim_{T \rightarrow \infty} \frac{1}{T} \ln \langle e^{TkA_T} \rangle, \quad k \in \mathbb{R}$$

Theorem: Gärtner (1977), Ellis (1984)

If $\lambda(k)$ is differentiable, then

- 1 LDP:

$$P(A_T = a) \approx e^{-T I(a)}$$

- 2 Rate function:

$$I(a) = \max_k \{ka - \lambda(k)\}$$

- $I(a)$ = Legendre transform of $\lambda(k)$

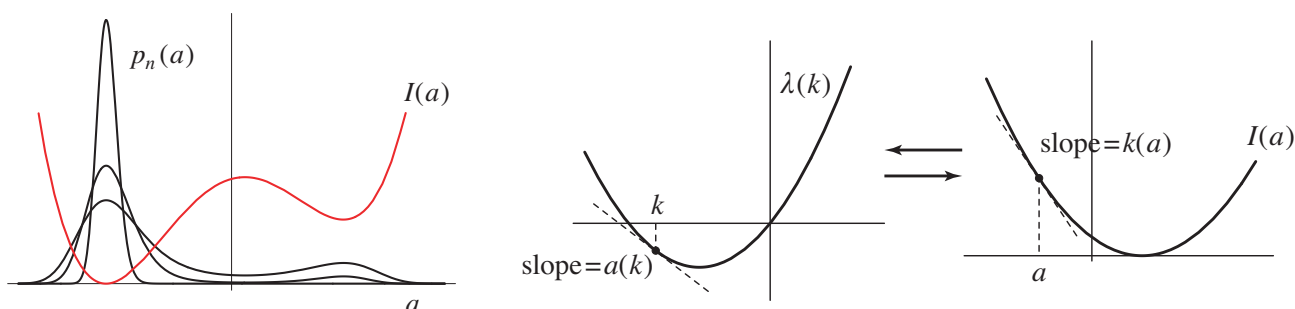


Richard S. Ellis



J. Gärtner

Important properties



- Duality:

$$\lambda'(k) = a \Leftrightarrow I'(a) = k$$

- Mean:

$$\lambda'(0) = \mu \Leftrightarrow I(\mu) = 0$$

- Variance:

$$\text{var}(A_T) \sim \frac{\lambda''(0)}{T}$$

- CLT:

$$I(a) = \frac{(a - \mu)^2}{2\sigma^2} + O(a^3)$$

Examples

- Sample mean:

$$S_n = \frac{1}{n} \sum_{i=1}^n X_i, \quad X_i \sim p(x), \quad \text{IID}$$

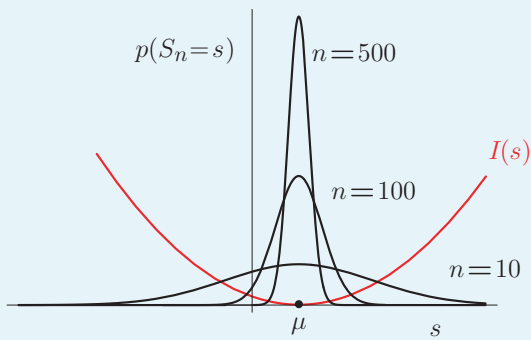
- SCGF:

$$\lambda(k) = \lim_{n \rightarrow \infty} \frac{1}{n} \ln E[e^{nkS_n}] = \ln E[e^{kX}]$$

Gaussian

$$\lambda(k) = \mu k + \frac{\sigma^2}{2} k^2, \quad k \in \mathbb{R}$$

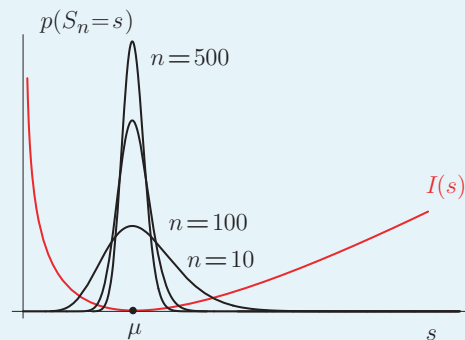
$$I(s) = \frac{1}{2\sigma^2} (s - \mu)^2, \quad s \in \mathbb{R}$$



Exponential

$$\lambda(k) = -\ln(1 - \mu k), \quad k < \frac{1}{\mu}$$

$$I(s) = \frac{s}{\mu} - 1 - \ln \frac{s}{\mu}, \quad s > 0$$



Large deviations of Markov processes

- Process:

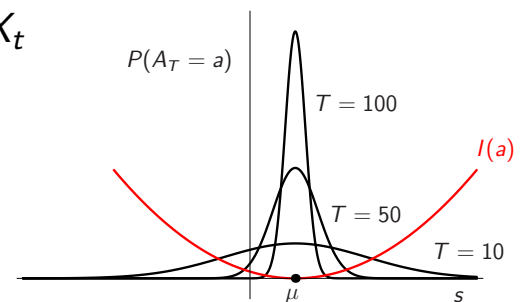
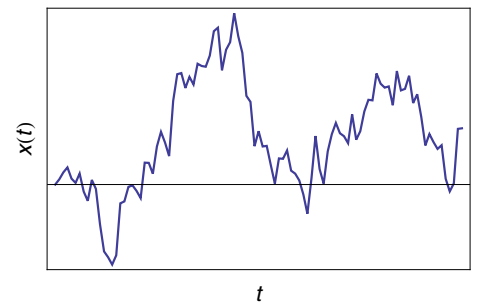
$$dX_t = F(X_t)dt + \sigma dW_t$$

- Observable:

$$A_T = \frac{1}{T} \int_0^T f(X_t) dt + \frac{1}{T} \int_0^T g(X_t) \circ dX_t$$

- LDP:

$$P(A_T = a) \approx e^{-T I(a)}, \quad T \rightarrow \infty$$



Question 6

Why study fluctuations?

Question 7

Do all observables satisfy the LDP?

Dual problem

Scaled cumulant function

$$\lambda(k) = \lim_{T \rightarrow \infty} \frac{1}{T} \ln E[e^{TkA_T}]$$

Gärtner-Ellis Theorem

$\lambda(k)$ differentiable, then

- 1 LDP for A_T
- 2 $I(a) = \sup_k \{ka - \lambda(k)\}$

Feynman-Kac-Perron-Frobenius

For Markov processes,

$$\lambda(k) = \zeta_{\max}(\mathcal{L}_k)$$

- SCGF = dominant eigenvalue
- Tilted (twisted) operator:

$$\mathcal{L}_k = F \cdot (\nabla + k\mathbf{g}) + \frac{1}{2}(\nabla + k\mathbf{g}) \cdot D(\nabla + k\mathbf{g}) + kf$$

- $\mathcal{L}_{k=0} = L$

Spectral problem

\mathcal{L}_k non-Hermitian in general

- Spectral problem:

$$\mathcal{L}_k r_k(x) = \lambda(k) r_k(x)$$

- Dual problem:

$$\mathcal{L}_k^\dagger l_k(x) = \lambda(k) l_k(x)$$

- Boundary condition:

$$r_k(x) l_k(x) \xrightarrow{|x| \rightarrow \infty} 0$$

- Normalization:

$$\int r_k(x) l_k(x) dx = 1, \quad \int l_k(x) dx = 1$$

- Extract dominant eigenvalue

Question 8

Why must wavefunctions decay at infinity?

Feynman-Kac formula

- Generating function:

$$G(x, t) = \left\langle e^{\int_0^t c(X_s) ds} \right\rangle_x = \langle e^{tkA_t} \rangle_x$$

- FK equation:

$$\partial_t G(x, t) = \mathcal{L}_c G(x, t), \quad \mathcal{L}_c = L + c$$

- Initial condition: $G(x, 0) = 1$

- Solution:

$$G(x, t) = (e^{t\mathcal{L}_k} 1)(x)$$

- Spectral decomposition:

$$G(x, t) = \sum_i e^{\zeta_i t} r_k^{(i)}(x) \sim \text{const} \cdot e^{\zeta_{\max} t}$$

Dominant eigenvalue is real (Perron-Frobenius)

Spectral problem

Equilibrium

- X_t reversible, g gradient, f arbitrary
- \mathcal{L}_k non-Hermitian but conjugated to Hermitian
- Real spectrum (quantum problem)
- Easy

Nonequilibrium

- X_t nonreversible OR g non-gradient, f arbitrary
- \mathcal{L}_k non-Hermitian, not conjugated to Hermitian
- Complex spectrum
- Full spectral problem
- Not so easy

Symmetrization

- Symmetrization:

$$\mathcal{H}_k = p_s^{1/2} \mathcal{L}_k p_s^{-1/2} = \sqrt{p_s} \mathcal{L}_k \frac{1}{\sqrt{p_s}}$$

- Hamiltonian:

$$\mathcal{H}_k = \frac{\varepsilon^2}{2} \Delta - V_k$$

- Effective potential:

$$V_k(x) = \frac{|\nabla U(x)|^2}{2\varepsilon^2} - \frac{\Delta U(x)}{2} - kf(x)$$

- Spectral problem:

$$\mathcal{H}_k \psi_k = \lambda(k) \psi_k,$$

- Eigenfunctions:

$$\psi_k(x) = \sqrt{p_s(x)} r_k(x), \quad \psi_k(x) = l_k(x) / \sqrt{p_s(x)}$$

- Normalization:

$$\psi(x)^2 \xrightarrow{|x| \rightarrow \infty} 0, \quad \int \psi(x)^2 dx = 1$$

Example

- Process:

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

- Observable:

$$A_T = \frac{1}{T} \int_0^T X_t dt$$

- Tilted generator:

$$\mathcal{L}_k = L + kf = -\gamma x \frac{d}{dx} + \frac{\sigma^2}{2} \frac{d^2}{dx^2} + kx, \quad k \in \mathbb{R}$$

- Stationary distribution:

$$p_s(x) = \sqrt{\frac{\gamma}{\pi \varepsilon^2}} e^{-\gamma x^2 / \varepsilon^2} \propto e^{-2U(x) / \varepsilon^2}$$

- Effective hamiltonian:

$$\mathcal{H}_k = \frac{\varepsilon^2}{2} \frac{d^2}{dx^2} - \frac{\gamma^2 x^2}{2\varepsilon^2} + \frac{\gamma}{2} + kx.$$

Example (cont'd)

- Dominant eigenvalue:

$$\lambda(k) = \frac{\varepsilon^2 k^2}{2\gamma^2}$$

- Rate function:

$$I(a) = \frac{\gamma^2 a^2}{2\varepsilon^2}$$

- Eigenfunctions:

$$\psi_k(x) = \left(\frac{\gamma}{\pi\varepsilon^2}\right)^{1/4} \exp\left(-\frac{\gamma(x - \varepsilon^2 k/\gamma^2)^2}{2\varepsilon^2}\right)$$

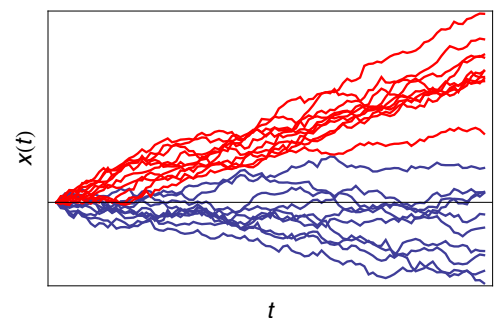
$$r_k(x) = \exp\left(\frac{kx}{\gamma} - \frac{3\varepsilon^2 k^2}{4\gamma^3}\right)$$

$$l_k(x) = \sqrt{\frac{\gamma}{\pi\varepsilon^2}} \exp\left(-\frac{\gamma(2x - \varepsilon^2 k/\gamma^2)^2}{4\varepsilon^2}\right)$$

How are fluctuations created?

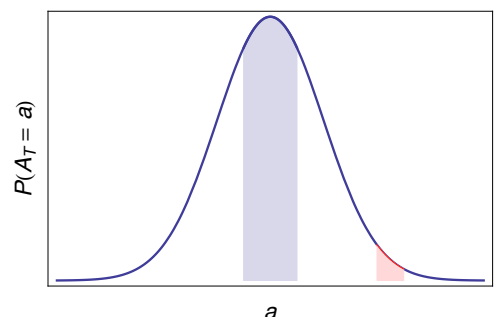
Physical problem

- Stochastic process: X_t
- Observable: A_T
- Look at trajectories leading to $A_T = a$
- Find effective process describing these trajectories



Mathematical problem

- Markov process: $\{X_t\}_{t=0}^T$
- Conditioned process: $X_t | A_T = a$
- Is it a Markov process?
- Construct its generator



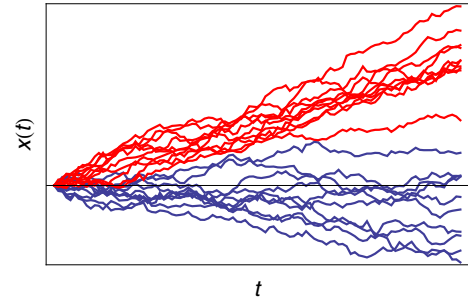
Fluctuation process

Driven diffusion

$$d\hat{X}_t = F_k(\hat{X}_t)dt + \sigma dW_t$$

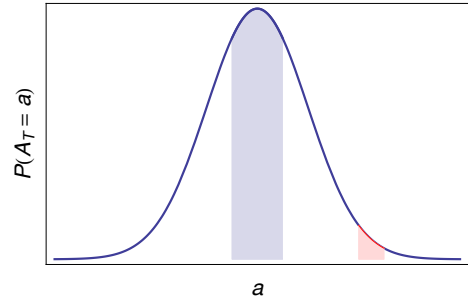
- Modified drift:

$$F_k(x) = F(x) + D(kg + \nabla \ln r_k), \quad I'(a) = k$$



Conditioning

$$\underbrace{X_t | A_T = a}_{\text{conditioned microcanonical}} \stackrel{T \rightarrow \infty}{\cong} \underbrace{\hat{X}_t}_{\text{driven canonical}}$$



Effective process creating the fluctuation

[Chetrite HT, PRL 2013, AHP 2015, JSTAT 2015]

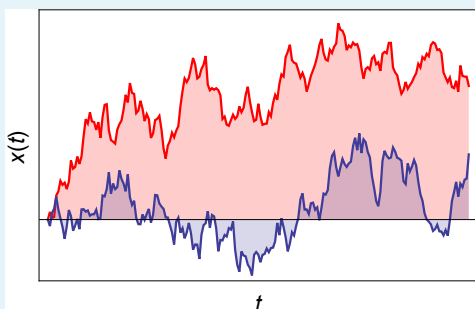
Examples

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

Area

$$A_T = \frac{1}{T} \int_0^T X_t dt = a$$

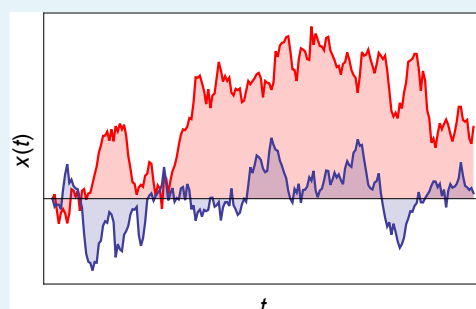
- $F_a(x) = -\gamma x + \gamma a$
- Modified drift



Empirical variance

$$A_T = \frac{1}{T} \int_0^T X_t^2 dt = a$$

- $F_a = -\sigma^2 x / (2a)$
- Modified friction



Research topics

- Jump processes
- Many particles dynamics
- Interacting particle systems
- Interacting diffusions
- Macroscopic fluctuation theory (stochastic PDEs)
- Fluctuation relations
- Level 2.5 large deviations
- Stochastic thermodynamics
- Entropy production
- Disordered systems (spin glasses, random graphs)
- Quantum systems (thermalization, nonequilibrium)
- Numerical large deviations
- Large deviation simulations

Final question

Does statistical mechanics have any relevance/future?

References



H. Touchette

The large deviation approach to statistical mechanics
[Physics Reports 478, 1-69, 2009](#)



H. Touchette

Introduction to large deviations: Theory, applications, simulations
[2011 Oldenburg School Lecture Notes, arxiv:1106.4146](#)



H. Touchette

Introduction to dynamical large deviations of Markov processes
[2017 FPSP Lecture Notes, arxiv:1705.06492](#)



More at

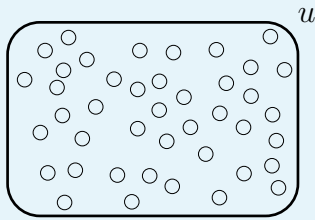
www.physics.sun.ac.za/~htouchette/ldt

www.physics.sun.ac.za/~htouchette/ldtcourse/

Extra: Analogy with equilibrium statistical mechanics

Microcanonical

[Einstein (1910)]



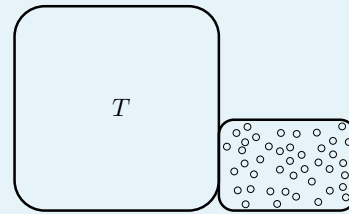
$$P_u(M_N = m) = e^{S(u,m)/k_B}$$

- Extensivity: $S \sim N$
- LDP:

$$P_u(M_N = m) \approx e^{-NI_u(m)}$$

Canonical

[Landau (1937)]



$$P_\beta(M_N = m) = e^{-F(\beta,m)}$$

- Extensivity: $F \sim N$
- LDP:

$$P_\beta(M_N = m) \approx e^{-NI_\beta(m)}$$

- Exponential concentration of probability
- Equilibrium states = minima and zeros of I