

# Poles of partition functions

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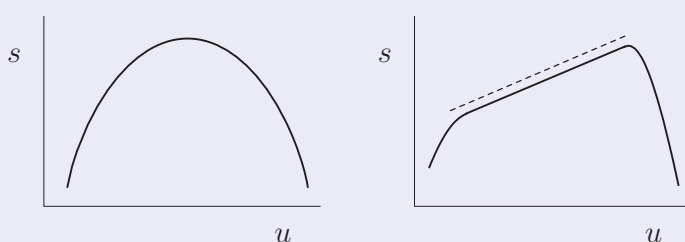
35th Conference of the Middle European Cooperation  
in Statistical Physics, Pont-à-Mousson, France  
March 2010

HT, Rosemary J. Harris (London), J. Tailleur (Edinburgh)  
arxiv:0912.3679  
Phys. Rev. E, Rapid Comm. 2010  
Slides at <http://www.maths.qmul.ac.uk/~ht/talks.html>

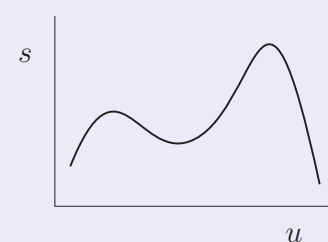
## Context

- Short-range vs long-range systems
  - ▶ Gravitating particles, unscreened plasmas, free electron laser  
[Campa, Dauxois, Ruffo, Physics Reports 2009]
- Nonconcave entropies
- Microcanonical vs canonical
- Nonequivalence of ensembles
- First-order phase transitions

### Short-range



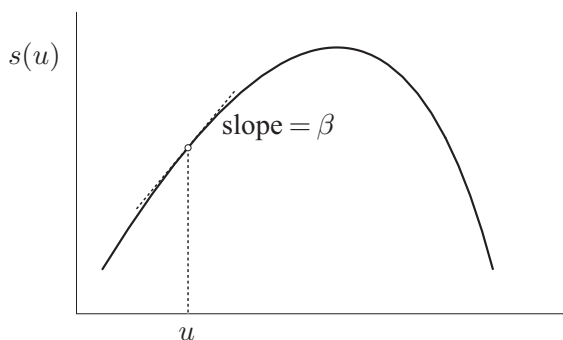
### Long-range



# Concave entropy

## Microcanonical

$$s(u) = \lim_{N \rightarrow \infty} \frac{1}{N} \ln \Omega_N(u)$$

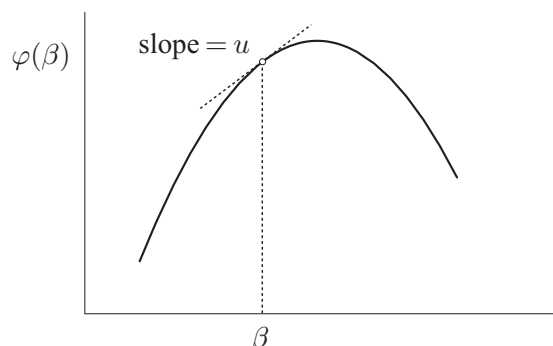


$$s(u) = \beta u - \varphi(\beta)$$

$$\varphi'(\beta) = u$$

## Canonical

$$\varphi(\beta) = \lim_{N \rightarrow \infty} -\frac{1}{N} \ln Z_N(\beta)$$



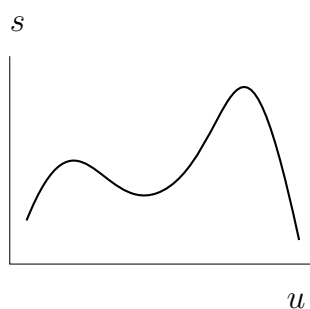
$$\varphi(\beta) = \beta u - s(u)$$

$$s'(u) = \beta$$

$$s \longleftrightarrow \varphi$$

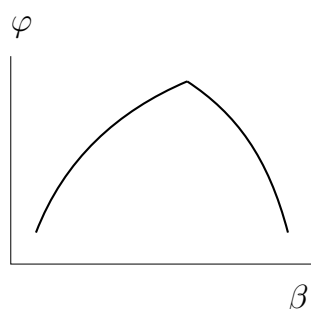
$$u \longleftrightarrow \beta$$

# Nonconcave entropies



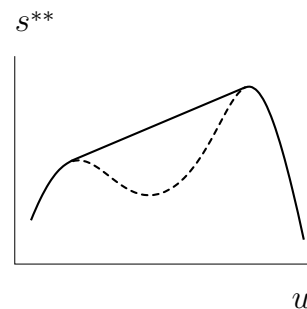
Nonconcave  
 $s$

$\xrightarrow{*}$



Always concave

$\xleftarrow{*}$   
 $\xrightarrow{*}$



$$s^{**} = \varphi^*$$

$$\varphi = s^*$$

$$s \neq s^{**}$$

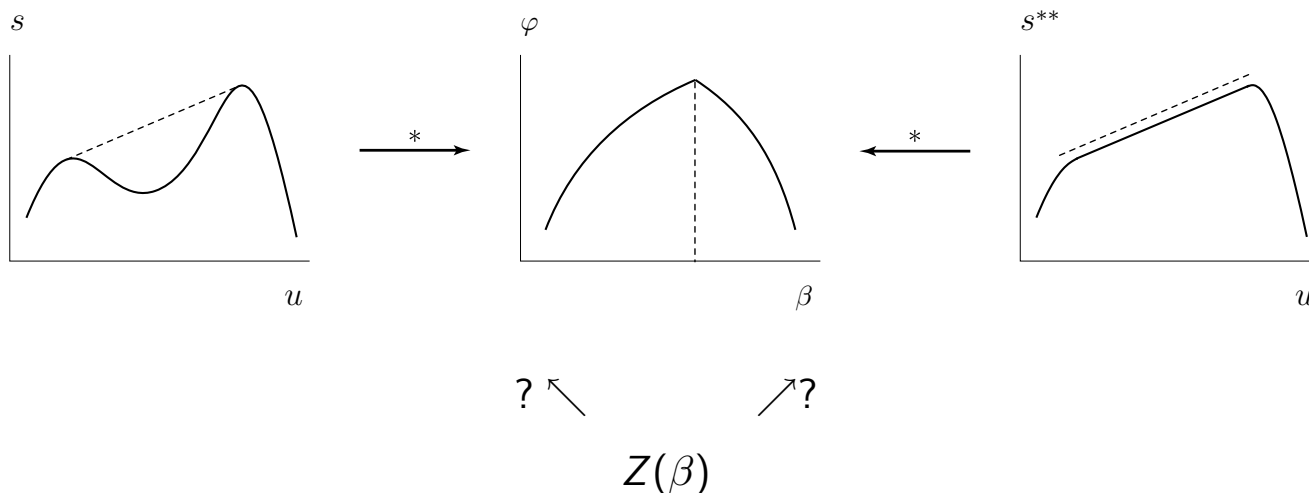
- $s^{**}(u) = \text{concave envelope of } s(u)$
- Nonequivalent ensembles
- Related to first-order phase transitions

# Problem

Microcanonical

Canonical

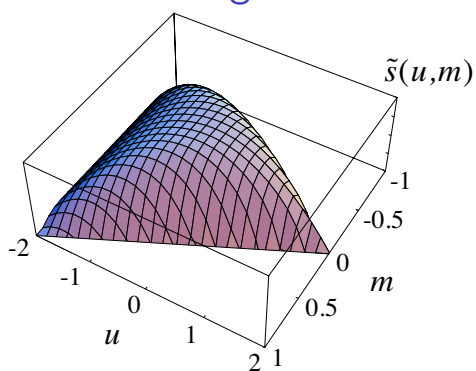
Microcanonical



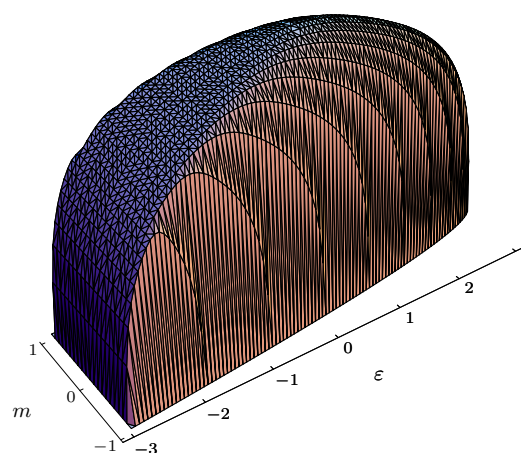
- Nonconcave/affine  $s(u)$  cannot be distinguished from  $\varphi(\beta)$
- Can they be distinguished from  $Z(\beta)$ ?

## Systems with affine entropies

2D Ising model



Spherical model



Kastner & Pleimling PRL 2009

- First-order phase transitions
- Metastability
- Phase separation
- Generic for short-range systems with first-order phase transition

## Result: Set-up

### Inverse Laplace transform

$$\Omega(u) = \frac{1}{2\pi i} \int_{r-i\infty}^{r+i\infty} Z(\beta) e^{\beta Nu} d\beta$$

Affine parts of  $s(u) \longleftrightarrow$  Poles in series representation of  $Z(\beta)$

### Ansatz

$$Z(\beta) = \sum_j c_j(\beta) e^{-N\varphi_j(\beta)}$$

- $\varphi_j(\beta)$  are independent of  $N$
- $c_j(\beta)$  are sub-exponential in  $N$
- $c_j(\beta)$  may have simple poles or branch cuts
- **Decomposition is not unique**

## Result

$$s(u) = \sup_j \sup_{\beta \in \{\beta_j^*, \beta_j^\times\}} \{\beta u - \varphi_j(\beta)\}$$

- $\beta_j^*$  = saddlepoints of  $\beta u - \varphi_j(\beta)$
- $\beta_j^\times$  = poles of  $c_j(\beta)$

### Affine $s(u)$

- poles is picked up
- constant saddlepoint is picked up (branch cut)

### Strictly concave or nonconcave $s(u)$

- no pole
- no constant saddlepoint

# Application 1

$$Z_1(\beta) = e^{-N\beta} + e^{N\beta}$$

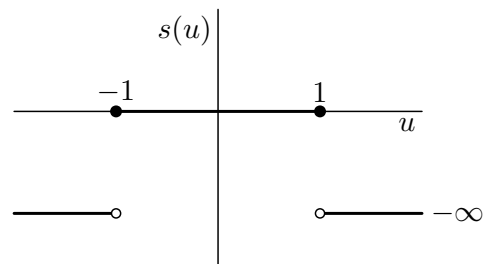
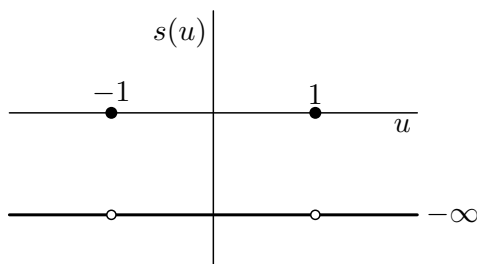
$$Z_2(\beta) = \frac{e^{N\beta} - e^{-N\beta}}{\beta}$$

$$\varphi(\beta) = -|\beta|$$

$$\varphi(\beta) = -|\beta|$$

$$\Omega(u) = \delta(u + 1) + \delta(u - 1)$$

$$\Omega(u) = \begin{cases} 1 & u \in [-1, 1] \\ 0 & \text{otherwise} \end{cases}$$

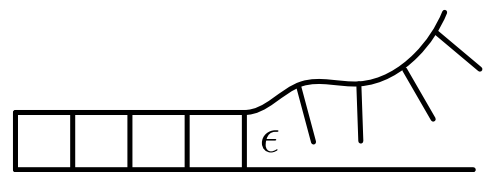


# Application 2: Kittel's DNA zipper model

Kittel Am. J. Phys. 1965

- Model:

- ▶  $N$  bonds
- ▶ Bond energy =  $\epsilon$
- ▶ Degeneracy =  $G$



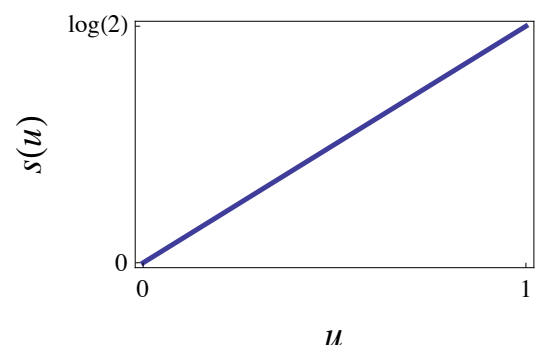
- Partition function:

$$Z(\beta) = \sum_{p=0}^{N-1} G^p e^{-\beta p \epsilon} \rightarrow \frac{1 - e^{-N(\beta \epsilon - \ln G)}}{N(\beta \epsilon - \ln G)}$$

- Pole:  $\beta_c = \epsilon^{-1} \ln G$

- Entropy:

$$s(u) = \begin{cases} \beta_c u & u \in [0, \epsilon) \\ -\infty & \text{otherwise} \end{cases}$$



# Conclusion

HT, R. J. Harris, J. Tailleur, arxiv:0912.3679

## Main result

- New mechanism for affine entropies
- General canonical calculation method for  $s(u)$
- Works for **affine** or **concave** or **nonconcave**  $s(u)$

## Applications

- Uhlenbeck-Kac model (1D gas)
- DNA models
- Markov processes (large deviation theory) [HT Physics Reports 2009]

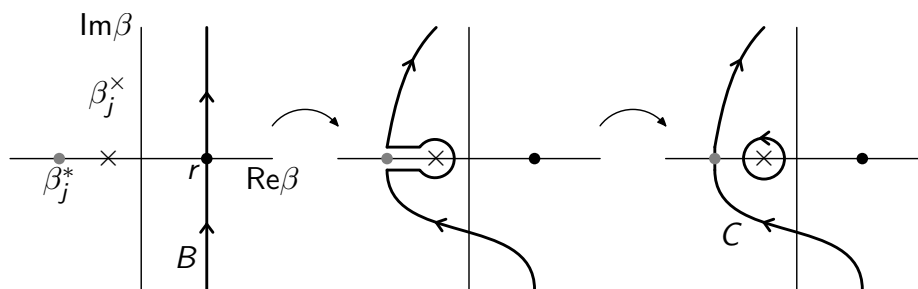
## Open problems

- $d > 1$  systems
- Real vs complex poles
- Other singularities?

## Idea of the proof

Deform Bromwich contour to S-D contour:

$$\frac{1}{2\pi i} \int_B c_j(\beta) e^{N[\beta u - \varphi_j(\beta)]} d\beta = \frac{1}{2\pi i} \int_C c_j(\beta) e^{N[\beta u - \varphi_j(\beta)]} d\beta + \sum \text{res}$$



Approximations:

$$\frac{1}{2\pi i} \int_B c_j(\beta) e^{N[\beta u - \varphi_j(\beta)]} d\beta \approx \underbrace{e^{N[\beta_j^* u - \varphi_j(\beta_j^*)]}}_{e^N \text{ S-D}} + \underbrace{\sum_{\ell} \sigma_{j\ell} e^{N[\beta_{j\ell}^x u - \varphi(\beta_{j\ell}^x)]}}_{e^N \text{ residue}}$$

- $\beta_j^*$ : Saddle-point
- $\beta_{j\ell}^x$ : Poles crossed (simple)
- $\sigma_{j\ell}$ : Residue parity (sign)

## Transfer matrix expansion

$$Z(\beta) = \text{Tr } T_\beta^{N-1} \rho_\beta = \sum_j c_j(\beta) e^{-N\varphi_j(\beta)}$$

- $T_\beta =$  Transfer matrix for  $e^{-\beta H}$
- $\rho_\beta =$  boundary condition vector
- $\varphi_j(\beta) = -\ln \xi_j(\beta)$
- $\xi_j(\beta) =$   $j$ th eigenvalue of  $T_\beta$
- $v_j(\beta) =$   $j$ th eigenvector
- $a_j(\beta) =$  projection of  $\rho_\beta$  onto  $v_j(\beta)$

$$c_j(\beta) = \frac{a_j(\beta)}{\xi_j(\beta)} \text{Tr } v_j(\beta)$$